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The diabatic contour-advective semi-Lagrangian algorithms for the shallow water equations on the sphere

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The extension of contour-advective semi-Lagrangian (CASL) algorithms to include diabatic forcing is discussed for the shallow water equations on the sphere. The CASL algorithms solve for potential vorticity (PV) by advecting particles in a contour representation, more precisely a piece-wise uniform, discontinuous representation. Such representation removes any need for differentiation/interpolation to advect PV and thus makes it possible to capture near-discontinuous distributions of PV generated by nonlinear advection. The CASL algorithms involve two novel transforms: (i) a contour-to-grid transform to find an Eulerian representation for the Lagrangian representation available in contours dividing piece-wise uniform regions of PV; (ii) a grid-to-contour transform to find a Lagrangian representation in the form of a set of piece-wise uniform regions of PV divided discontinuously by contours for the available Eulerian representation.

The diabatic CASL partitions the PV fields Q to an adiabatic part Q_a and a diabatic part Q_d . The adiabatic part Q_a is evolved in time by contour advection as in adiabatic CASL. No source/sink comes into the advection of Q_a , except a *regularisation* procedure called contour surgery. Having an Eulerian, grid representation only, the diabatic part Q_d is advected by a semi-Lagrangian advection scheme incorporating source/sink effects on PV. At certain time intervals, the Q_a and Q_d are merged on an Eulerian fine/ultra-fine grid and *recontoured* to reconstruct the Lagrangian representation. The recontouring is nothing but the realisation of the grid-to-contour transform. Ideally, we would like the merger process to be complete, i.e. reset Q_d to zero after recon-

touring. But in general the grid-to-contour transform leaves a non-zero *residual*. The treatment of residual and further developments on diabatic CASL algorithms, including a fully-Lagrangian description of PV dispensing with semi-Lagrangian advection, are discussed in detail.